# IMPERIAL

# **Nonlinear Dynamics of Active Filaments**

# Motivation

**Active filaments** are thin, elastic, slender rods that are commonly observed in nature. Driven by motor protein activity, active filaments play a crucial role in cellular processes such as fluid transport, mixing, and swimming, among many others.

Our goal is to analyse filament dynamics in biological settings using our numerical model, identify self-sustaining states (limit cycles), detect transitions between these states (bifurcations), and compute the stability of these states.

# The Model and Methods

- Molecular motors exert compressive forces on the filament and create a slip-flow over filament surface.
- Discretise the filament into self-locomoting segments.
- Solve the equations in the Stokes limit<sup>1</sup>: velocity and forces are linearly related!
- Model the squirming effect (effect of motor proteins) as a nonlinear forcing and solve for the velocity using Faxén's Laws.<sup>2</sup>

$$egin{bmatrix} m{V} \ m{\Omega} \end{bmatrix} = \mathcal{M}^{\mathsf{RPY}} egin{bmatrix} m{F}^H \ m{T}^H \end{bmatrix} + \mathcal{M}^{\mathsf{sq}} egin{bmatrix} m{H(b_1)} \ m{0} \end{bmatrix}$$

- $b_1$  measures motor protein activity.
- Analyse recurrent patterns in the dynamics using Poincaré maps.
- Use a Jacobian-free Newton-Krylov (JFNK) method to obtain accurate representations of periodic states and compute their Floquet stability.
- Track each solution branch to find the bifurcations.

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pproach employed to describe the motor protein activity and resulting filament dynamic proting squirmers. a is the radius of the filament's cross-section, L is the length of the filament,  $K_B$  is filament stiffness, and  $\nu$  is fluid viscosity

**Table 1:** Bifurcations of the squirming filament, with  $b_1$  values given in the bottom row. DH: double Hopf bifurcation, PB: pitchfo bifurcation. HB: Hopf bifurcation

Event	DH	PB	HB1	HB2	
$b_1$	19.1	80.8	90.8	184	

# Results

We identify five distinct states: upright, whirling, beating, QP1, and QP2. Upright is steady; whirling and beating are periodic orbits; QP1 and QP2 are quasi-periodic states.

### References

<sup>1</sup>Schoeller et al. JCP, 2020. <sup>2</sup>Clarke, B. PhD Thesis. <sup>3</sup>http://ladyofhats. com/ <sup>4</sup>https://www.sciencephoto.com/

Double Hopf bifurcation: a dynamic instability with two distinct eigenmodes. The phase difference between the modes determines the state: beating  $(\Delta \phi = 0)$ , and whirling  $(\Delta \phi = \frac{\pi}{2})$ 

Bistability of whirling and beating in the range  $b_1 \in [80.8, 90.8].$ 

 Modal decomposition of QP1 into whirling and beating:  $\boldsymbol{u}_{QP1}(t) \approx c_W \boldsymbol{u}_W(t) + c_B \boldsymbol{u}_B(t+\phi)$ , displaying both characteristics.

At high actuation, transition to QP2 – a coiled filament with changing winding direction.

• Robust dynamics to changes in motor protein activity  $b_1$ . QP2 does not transition to chaos!



#### **Future Work**

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t, whirling, and beating states: instability growth rate ( $\lambda$ ) vs. motor protein activity ( $b_1$ 

Data-driven analysis of the dynamics to explain the hydrodynamic interactions leading to the observed filament behaviours.

Extending the current framework to investigate collective dynamics of multiple filaments.

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