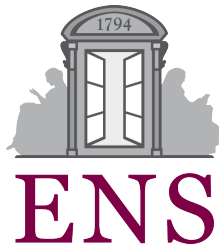


***Basic Facts of GFD +
Atmospheric LFV, Wind-driven Oceans,
Paleoclimate & “Tipping Points”***

Michael Ghil

Ecole Normale Supérieure, Paris, and
University of California, Los Angeles



Please visit these sites for more info.

<http://www.atmos.ucla.edu/tcd/>

<http://www.environnement.ens.fr/>

Overall Outline

- **Lecture I: Observations and planetary flow theory (GFD^(⌘))**
- **Lecture II: Atmospheric LFV^(*) & LRF^(**)**
- **Lecture III: EBMs⁽⁺⁾, paleoclimate & “tipping points”**
- **Lecture IV: The wind-driven ocean circulation**
- **Lecture V: Advanced spectral methods—SSA^(±) *et al.***
- **Lecture VI: Nonlinear & stochastic models—RDS^(◇)**

(⌘) GFD = Geophysical fluid dynamics

(*) LFV = Low-frequency variability

(**) LRF = Long-range forecasting

(+) EBM = Energy balance model

(±) SSA = Singular-spectrum analysis

(◇) RDS = Random dynamical system

Lecture I: Observations and Basic Planetary Flow Theory

Outline

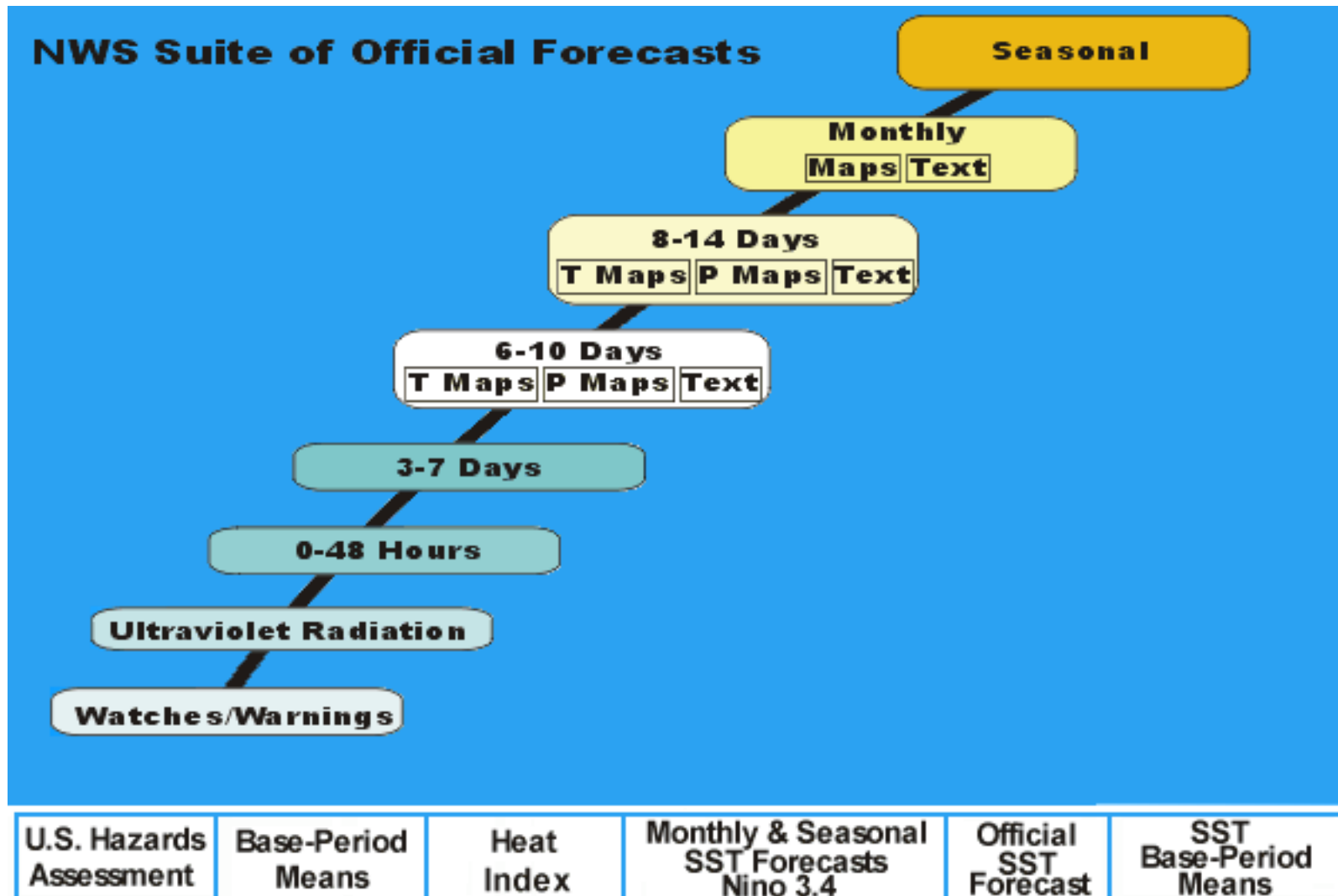
1. General introduction and motivation
 - Scale dependence of atmospheric & oceanic flows
 - Turbulence & predictability
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 - The atmospheric heat engine
 - Shallowness
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Weather & climate: variability and prediction, I

U.S. National Weather Service (NWS): Forecast suite



Weather & climate: variability and prediction

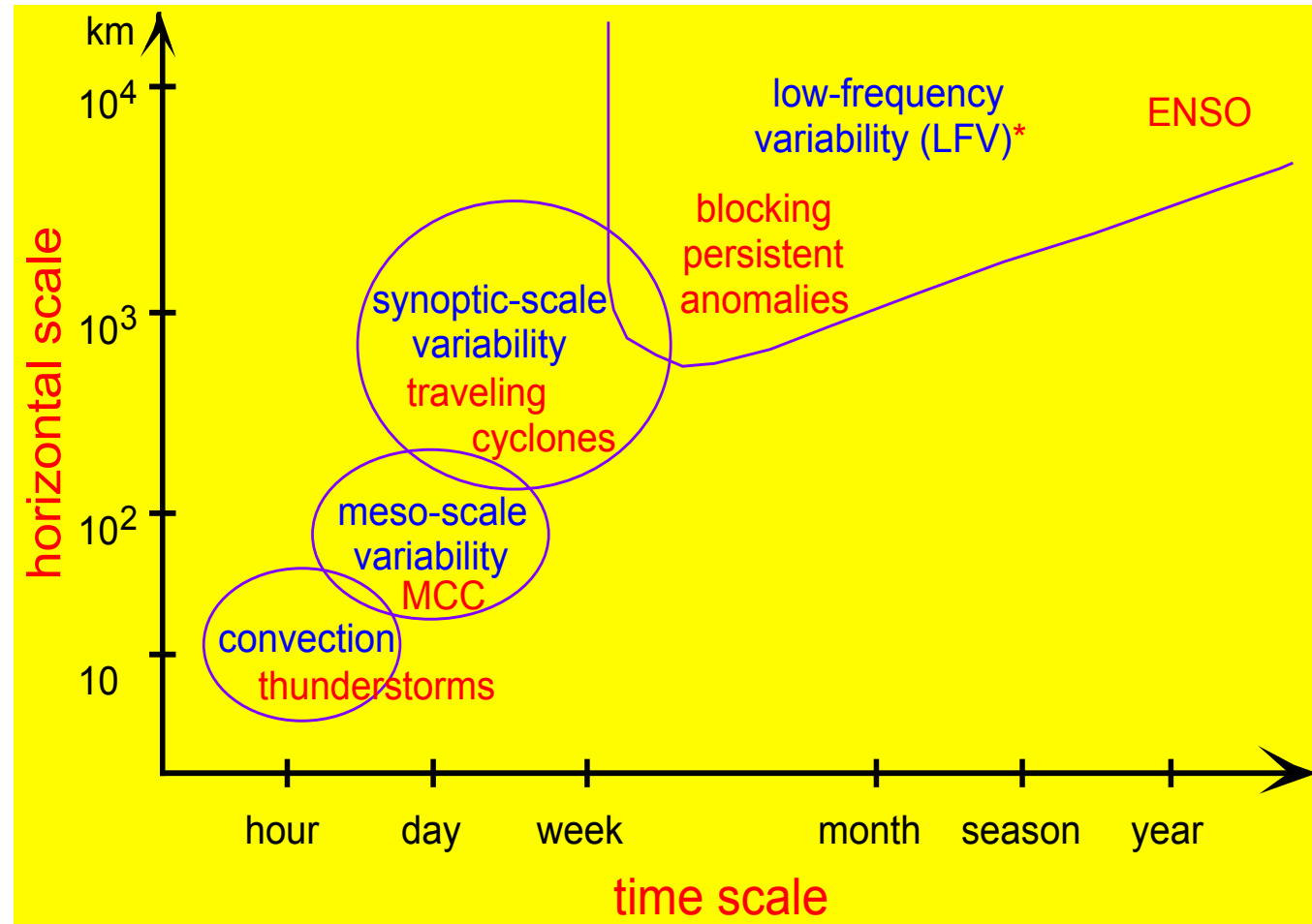
Problem 1: Find the comparable forecast suites on the web sites of the UK Met Office & the ECMWF

Weather & climate: Observations, II

Space & time scales, $k \sim \omega^{(*)}$

Atmospheric LFV \approx
10–100 days
(intraseasonal)

Oceanic LFV \approx
3–300 years
(interannual–
interdecadal)

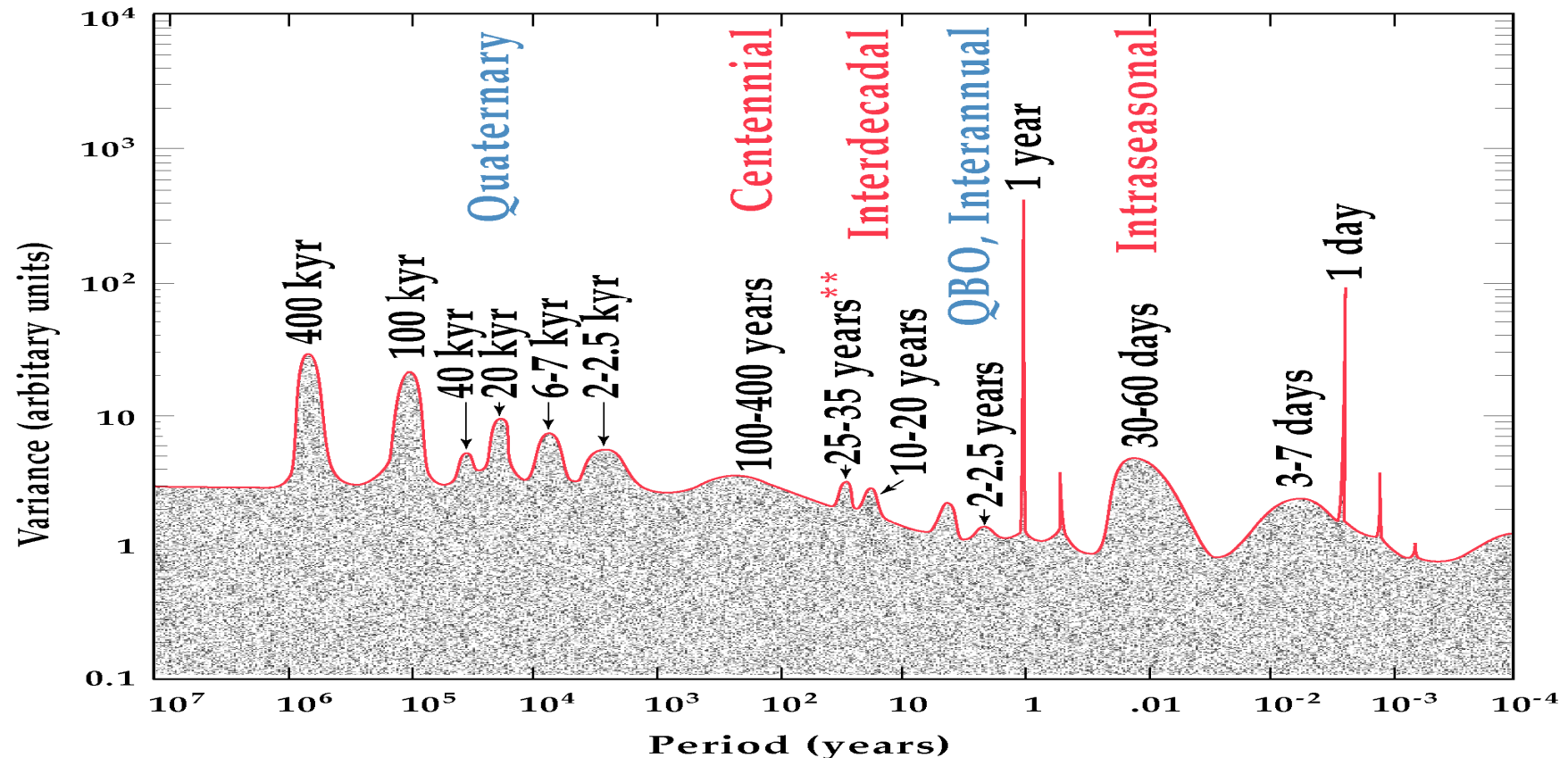


(*) A high-variability ridge lies close to the diagonal of the plot
(cf. also Fraedrich & Böttger, *JAS*, 1978)

Composite spectrum of climate variability

Standard treatment of frequency bands:

1. High frequencies – white noise (or “colored”)
2. Low frequencies – slow evolution of parameters

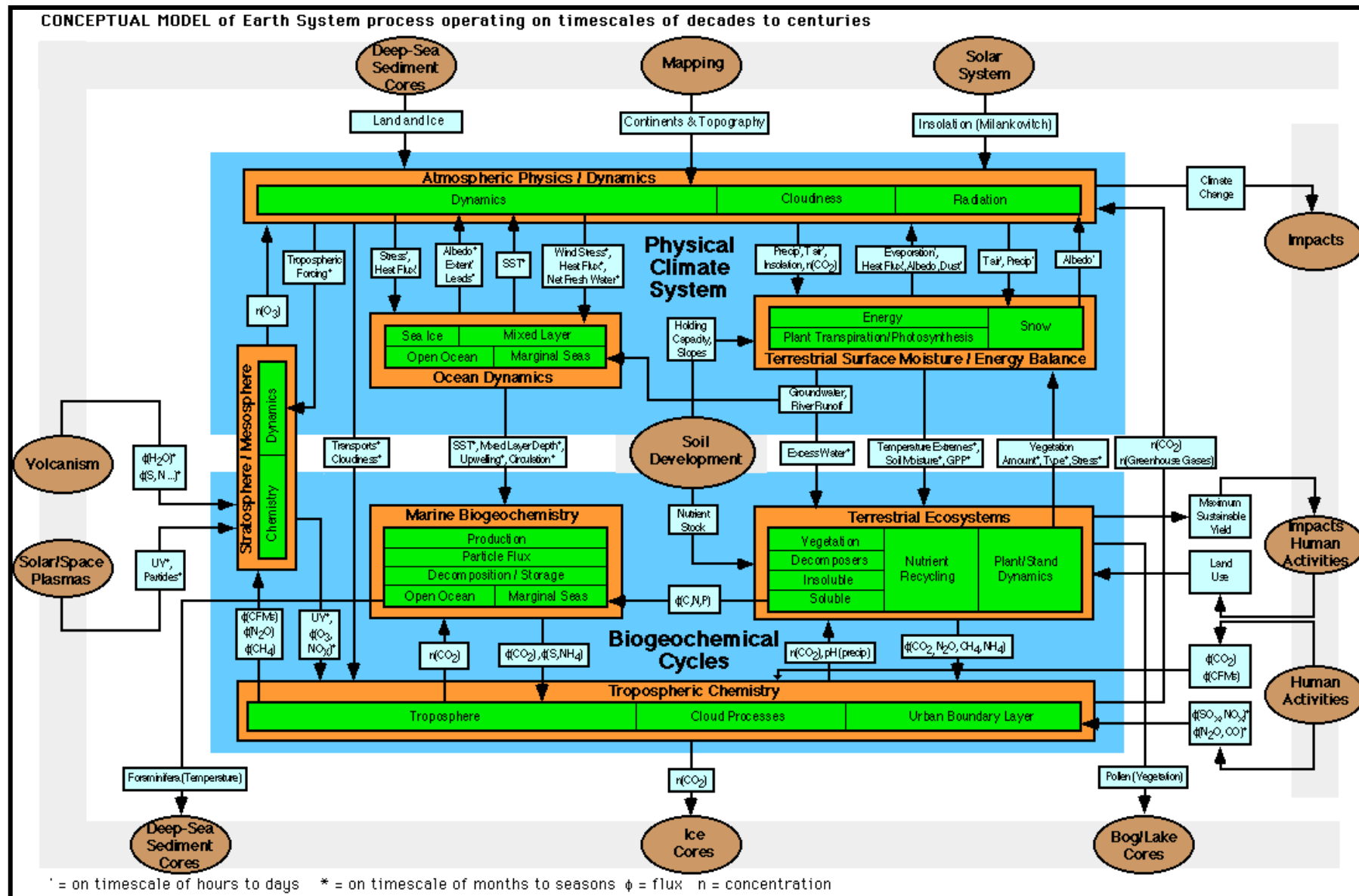


From Ghil (2001, *EGEC*), after Mitchell* (1976)

* “No known source of deterministic internal variability”

** 27 years – Brier (1968, *Rev. Geophys.*)

F. Bretherton's "horrendogram" of Earth System Science



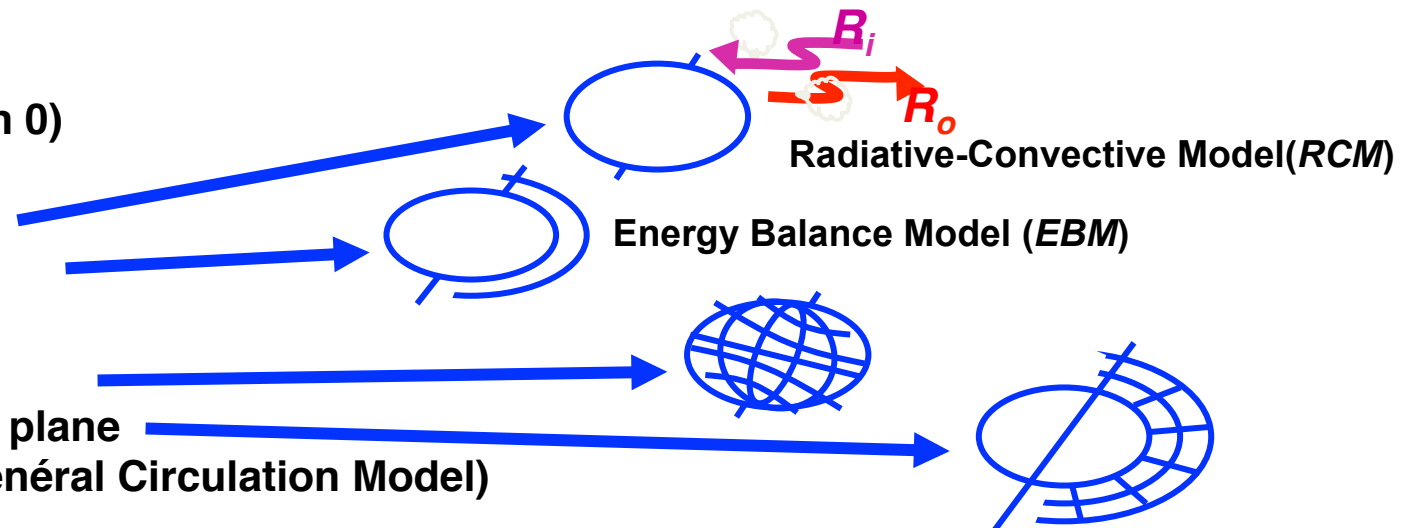
Climate models (atmospheric & coupled) : A classification

• Temporal

- stationary, (quasi-)equilibrium
- transient, climate variability

• Space

- 0-D (dimension 0)
- 1-D
 - vertical
 - latitudinal
- 2-D
 - horizontal
 - meridional plane
- 3-D, GCMs (Général Circulation Model)
 - horizontal
 - meridional plane
- Simple and intermediate 2-D & 3-D models



• Coupling

- Partial
 - unidirectional
 - asynchronous, hybrid
- Full

Hierarchy: from the simplest to the most elaborate,
iterative comparison with the observational data

ITALIAN PHYSICAL SOCIETY

PROCEEDINGS
OF THE
INTERNATIONAL SCHOOL OF PHYSICS
«ENRICO FERMI»

COURSE LXXXVIII

edited by M. GHIL
Director of the Course
and by R. BENZI and G. PARISI
VARENNA ON LAKE COMO
VILLA MONASTERO
14 - 24 June 1983

*Turbulence and Predictability
in Geophysical Fluid Dynamics
and Climate Dynamics*

1985



NORTH-HOLLAND
AMSTERDAM · OXFORD · NEW YORK · TOKYO

Lecture I: Outline

1. General introduction

- Scale dependence of atmospheric
- **Turbulence & predictability**

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The Lorenz model (1963a): a concrete example of a strange attractor^(*)

- *The model equations: 3 coupled, nonlinear ODEs*

$$\dot{X} = -\sigma X + \sigma Y \quad (1)$$

$$\dot{Y} = -XZ + rX - Y \quad (2)$$

$$\dot{Z} = XY - bZ \quad (3)$$

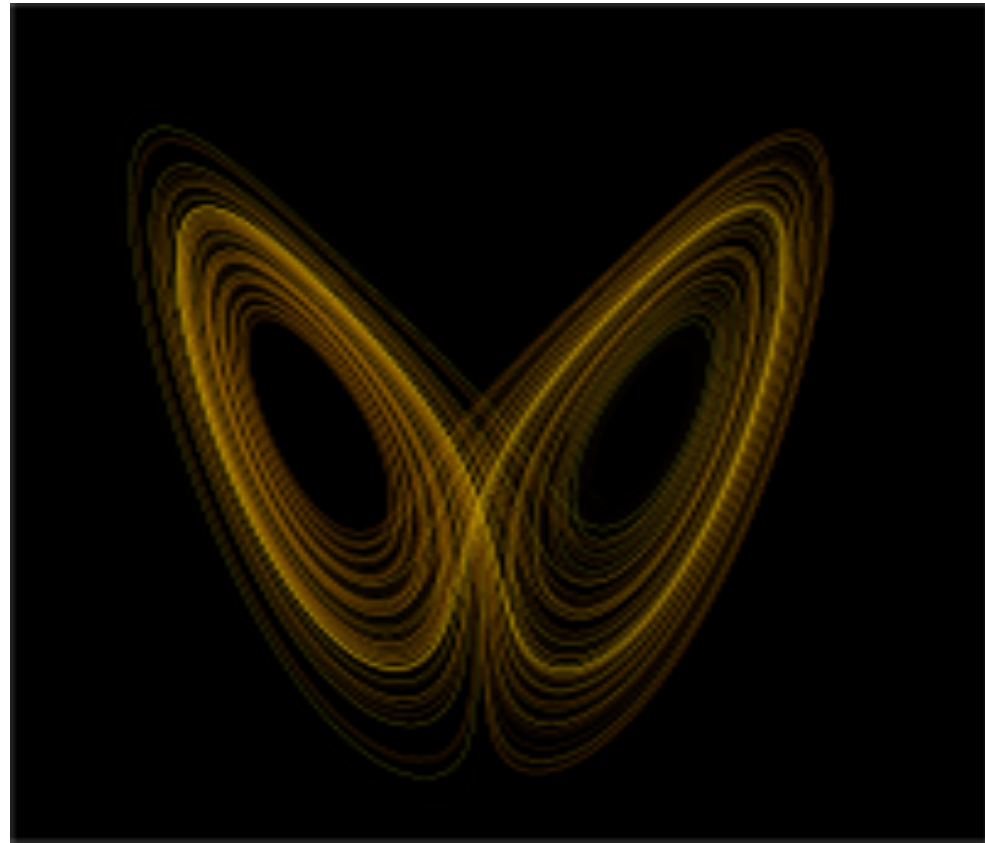
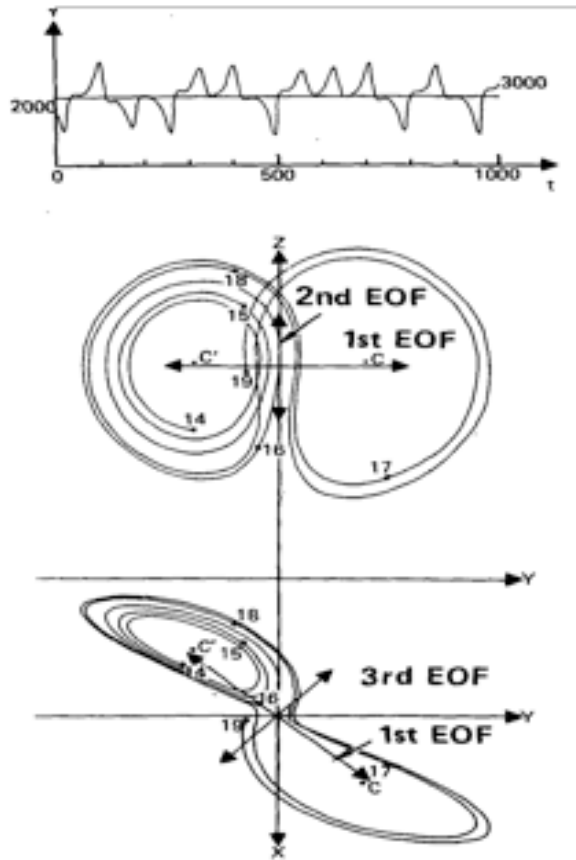
- **Physics: a model of thermal convection in 2-D**

The variables X and Y represent the intensity of the **velocity** field in a 2-D space, Z is the deviation of the vertical **temperature profile** from pure conduction (no motion), and $(X, Y, Z)^\bullet$ is their rate of change.

The parameters are the **Rayleigh** number ρ (intensity of the thermal forcing), the **Prandtl** number σ (the fluid's dissipative properties) and β characterizes the **wave length** of the perturbation from pure conduction.

^(*) Mommy, what's a strange attractor, please?

The Lorenz convection (1963a) model – some numerical solutions



Plot of $Y = Y(t)$ + projections
onto the (X, Y) & (Y, Z) planes

Trajectory in phase space

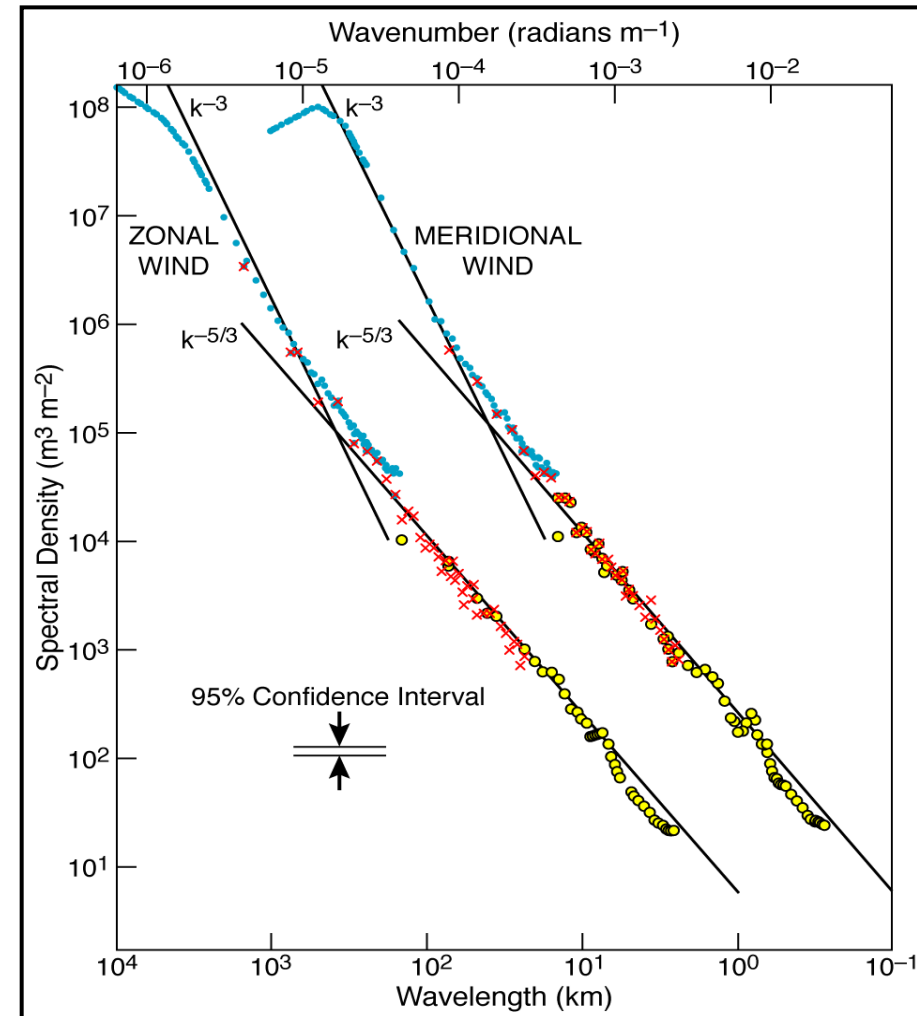
Both for the canonical “chaotic” values $\rho = 28$, $\sigma = 10$, $\beta = 8/3$.

The Lorenz (1963a) convection model

Problem 2: Find the appropriate software to compute the Lorenz “butterfly” and use it to do so.

But chaos doesn't explain everything: there are many other sources of irregularity!

- Indeed, the atmosphere's & oceans' energy spectrum is “full” – all the time & space scales are active, and contribute to prediction errors.
- Still, one can imagine that the longest, slowest & most energetic modes play a key role.
- “One person's signal is another person's noise.”



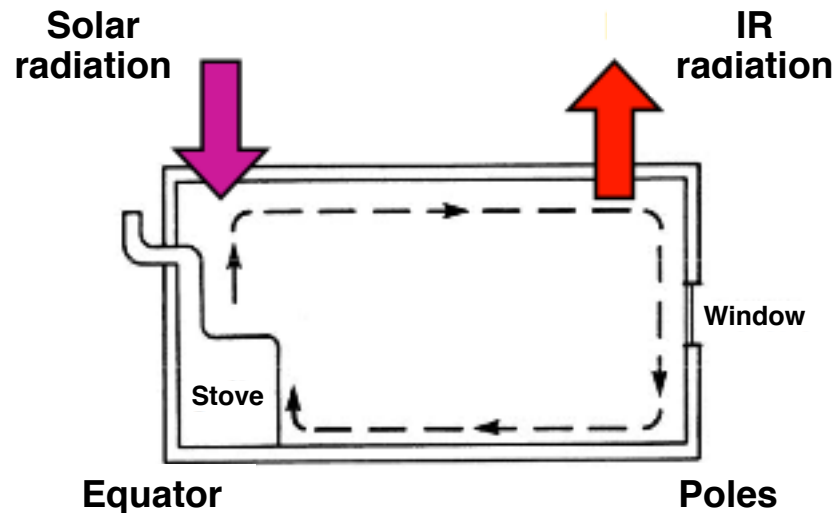
After Nastrom & Gage (JAS, 1985)

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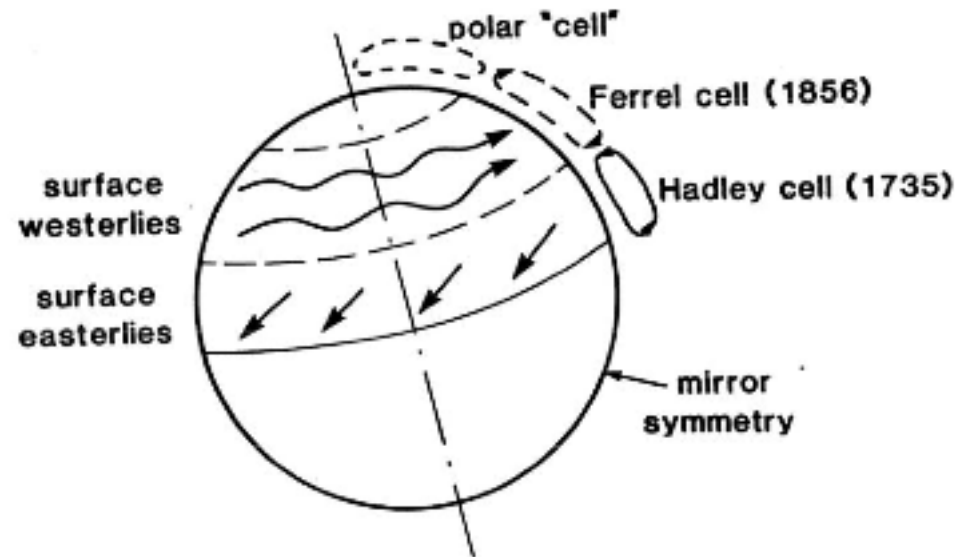
The mean atmospheric circulation

Direct Hadley circulation



Idealized view of the atmosphere's global circulation.*

Observed circulation



Schematic diagram of the atmospheric global circulation.*

*From Ghil and Childress (1987), Ch. 4

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Basic facts of large-scale atmospheric life, or how to read weather maps – I

1. Shalowness, I

$$\delta = H/L \ll 1 \quad \Rightarrow$$

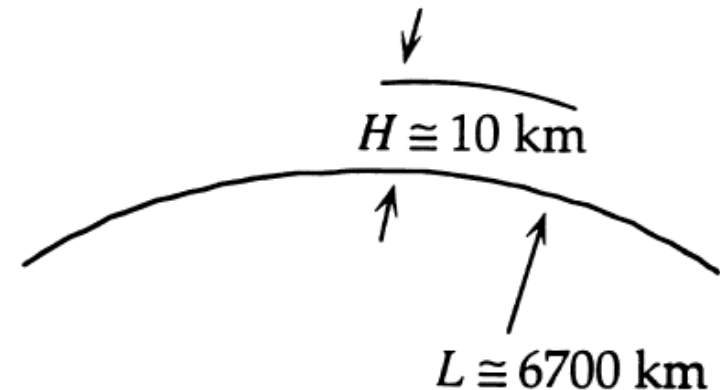
The flow is approximately 2-D (“barotropic”) & hence, to a good approximation, it is governed by shallow-water equations (SWE):

$$u_t + uu_x + vu_y = -\phi_x + fv,$$

$$v_t + uv_x + vv_y = -\phi_y - fu,$$

$$\phi_t + (u\phi)_x + (v\phi)_y = 0.$$

Here h is the height of the “free surface,” which is of order H , while $\phi = gh$ is the *geopotential*.



Basic facts of large-scale atmospheric life, or how to read weather maps - II

1. Shallowness, II

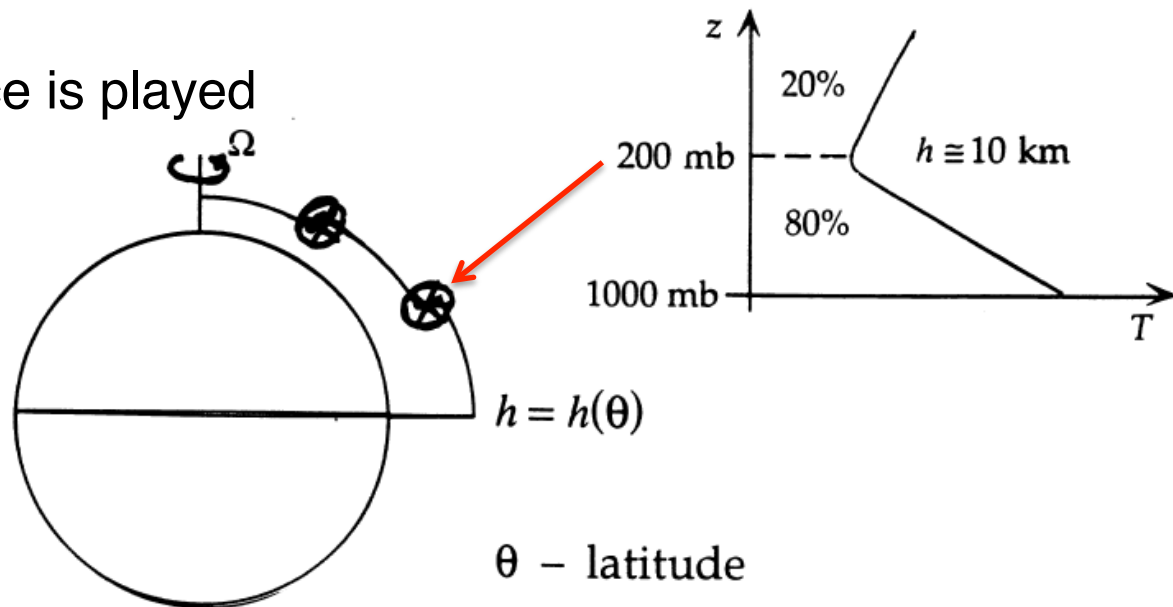
$\delta = H/L \ll 1$ also implies that the flow is approximately

hydrostatic, $p_z = -\rho g < 0$; hence “pressure coordinates”:

$$z_p = -\frac{1}{gp} \text{ or } \phi_p = -\frac{RT(p)}{p}.$$

The role of the free surface is played by the *tropopause*.

The atmospheric jets coincide roughly with the “tropopause gaps.”



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Basic facts of large-scale atmospheric life, or how to read weather maps – III

2. Rotation & geostrophy

$f = 2\Omega \sin \theta$ is the *planetary vorticity*, or the *Coriolis parameter*.

The Rossby number $\epsilon = U/fL$ measures the importance of rotation:

It's important if ϵ is small: $\epsilon \ll 1$.

In *geostrophic flow*, $\epsilon \rightarrow 0$ and thus the SWE are reduced to

$$u = -(1/f)\phi_y, \quad v = (1/f)\phi_x.$$

The flow is parallel to isobaric contours, rather than perpendicular, and thus

$\psi = (g/f)h$ is a stream function.

In the quasi-geostrophic approximation, $0 < \epsilon \ll 1$

allows for small, slow deviations from exact geostrophy.

Basic facts of large-scale atmospheric life – IV

3. *Rotation + shallowness* → The quasi-geostrophic, equivalent-barotropic potential vorticity equation with topography

$$(\Delta - \lambda^{-2})\eta_t + J(\eta, \Delta - \lambda^{-2}\eta + h_0) = 0;$$

here Δ is the *Laplacian*,
 $J(\eta, Q)$ is the *Jacobian*, $J(\eta, Q) = \frac{\partial \eta}{\partial x} \frac{\partial Q}{\partial y} - \frac{\partial \eta}{\partial y} \frac{\partial Q}{\partial x} = (u, v) \cdot \nabla Q$

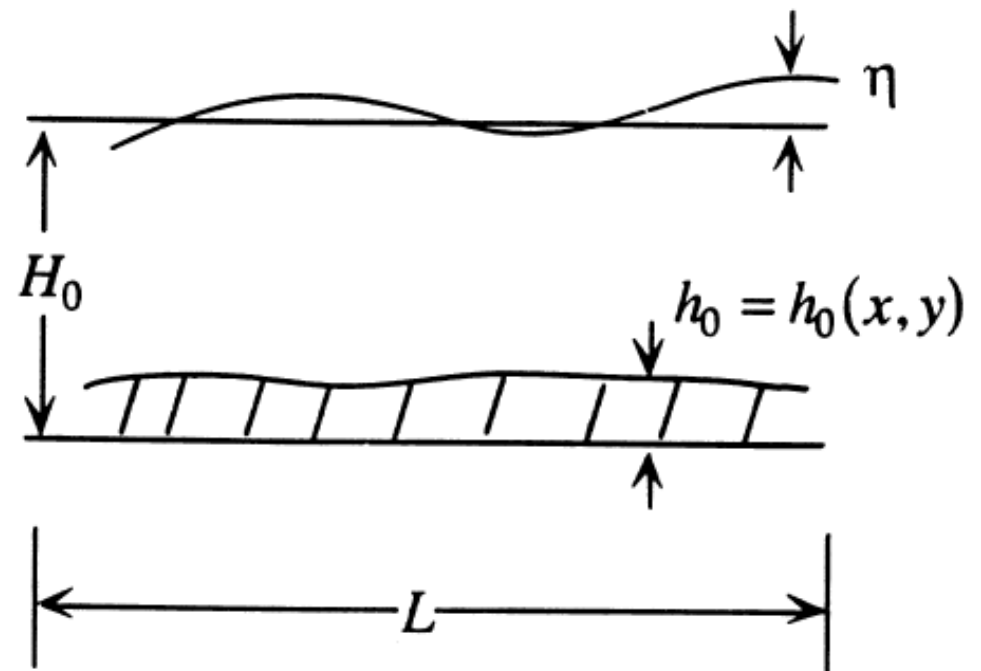
$$h = H_0(1 + \epsilon\lambda^{-2}\eta), \quad h_0 = H_0\epsilon h_0^*.$$

The *potential vorticity* Q equals

$$Q = (\Delta - \lambda^{-2})\eta + h_0,$$

and the *Rossby deformation radius* L_R plays a key role in it,

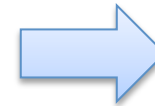
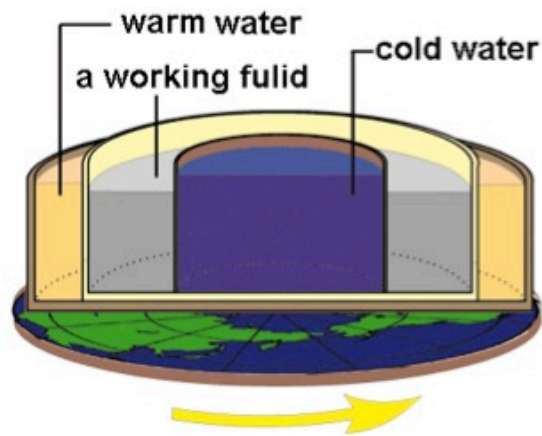
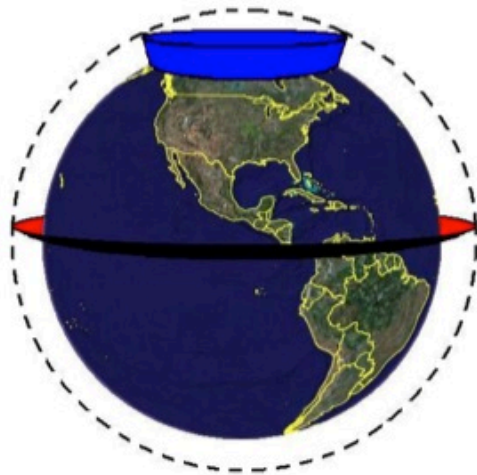
$$\lambda = L/L_R, \quad L_R = (gH_0)^{1/2}/f_0.$$



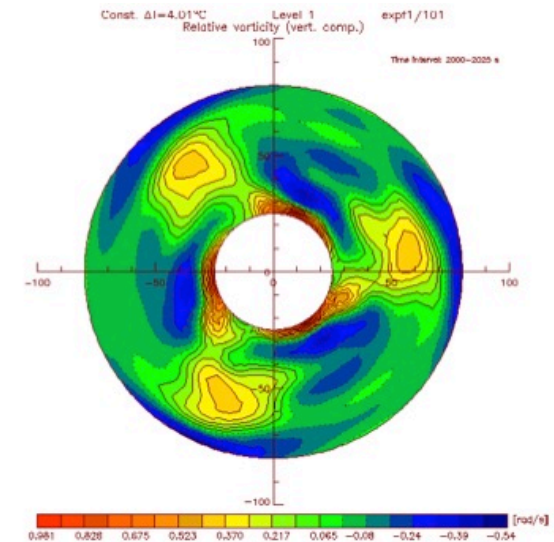
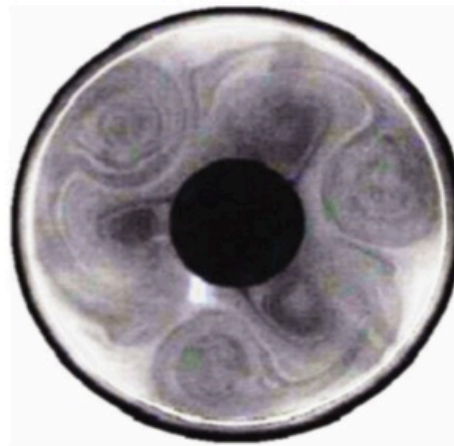
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Laboratory Analogues of Planetary Atmospheric Circulation Systems



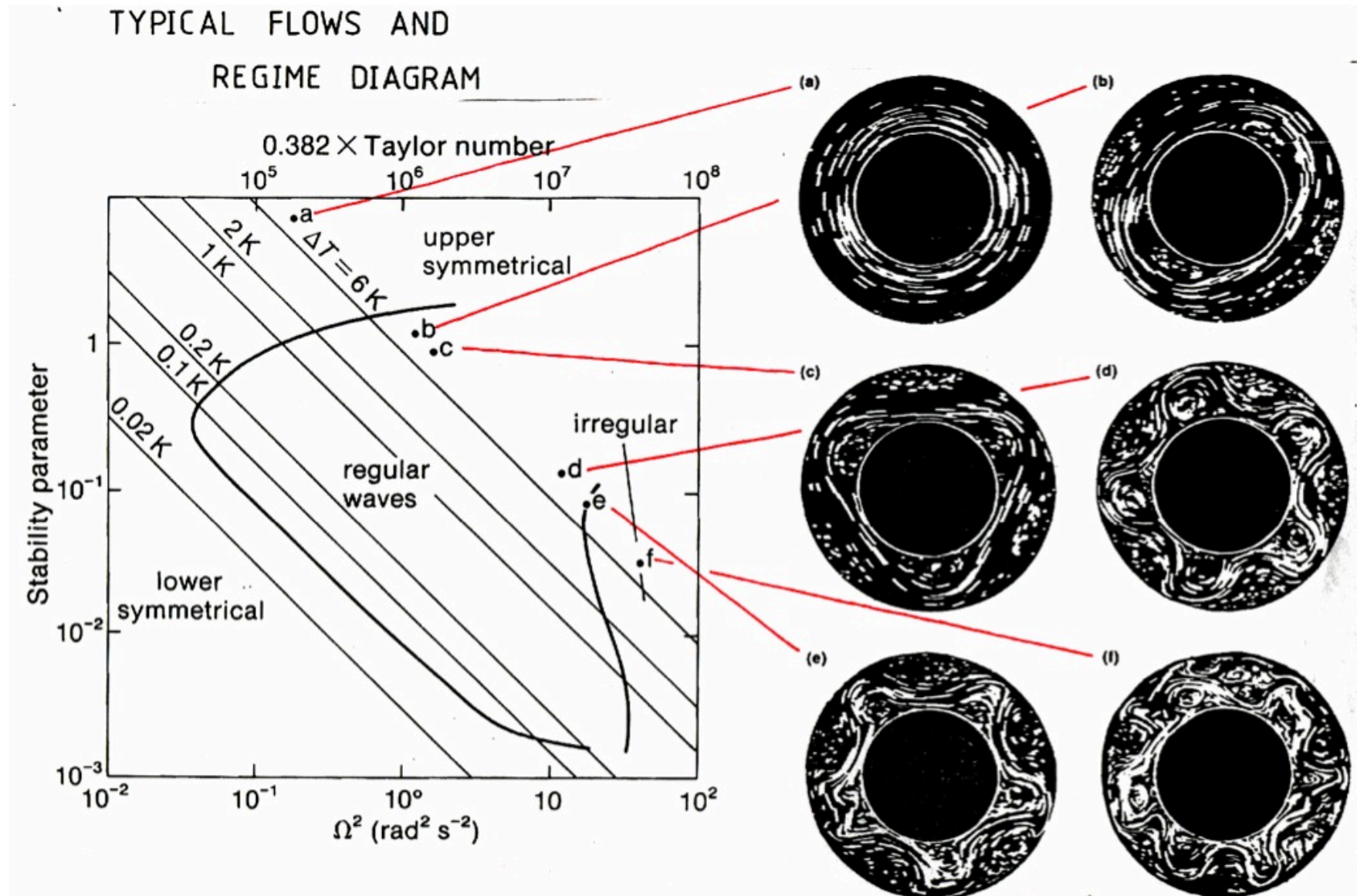
- *Baroclinic instability:*
 - A potential energy releasing instability in the atmosphere and the oceans



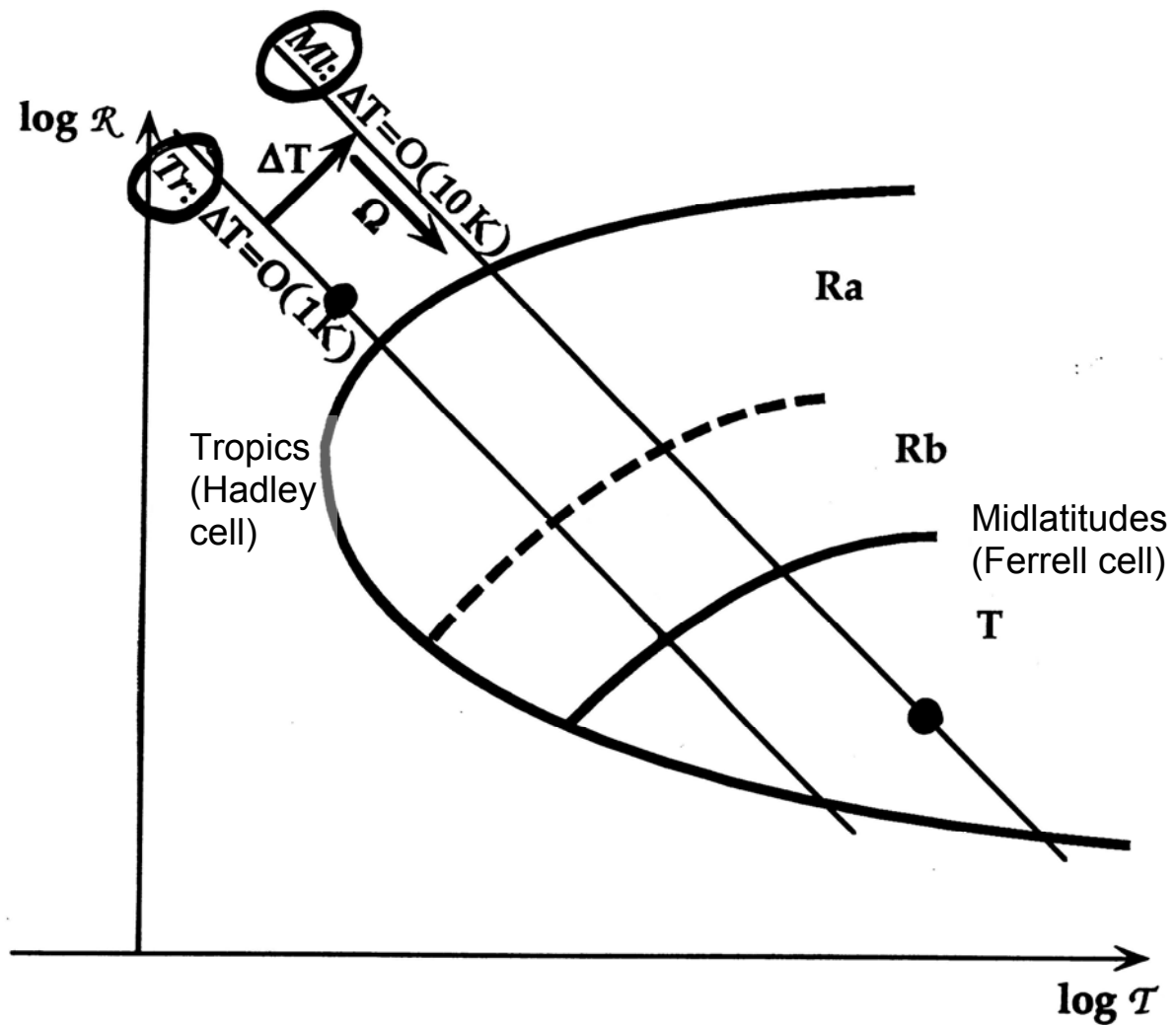
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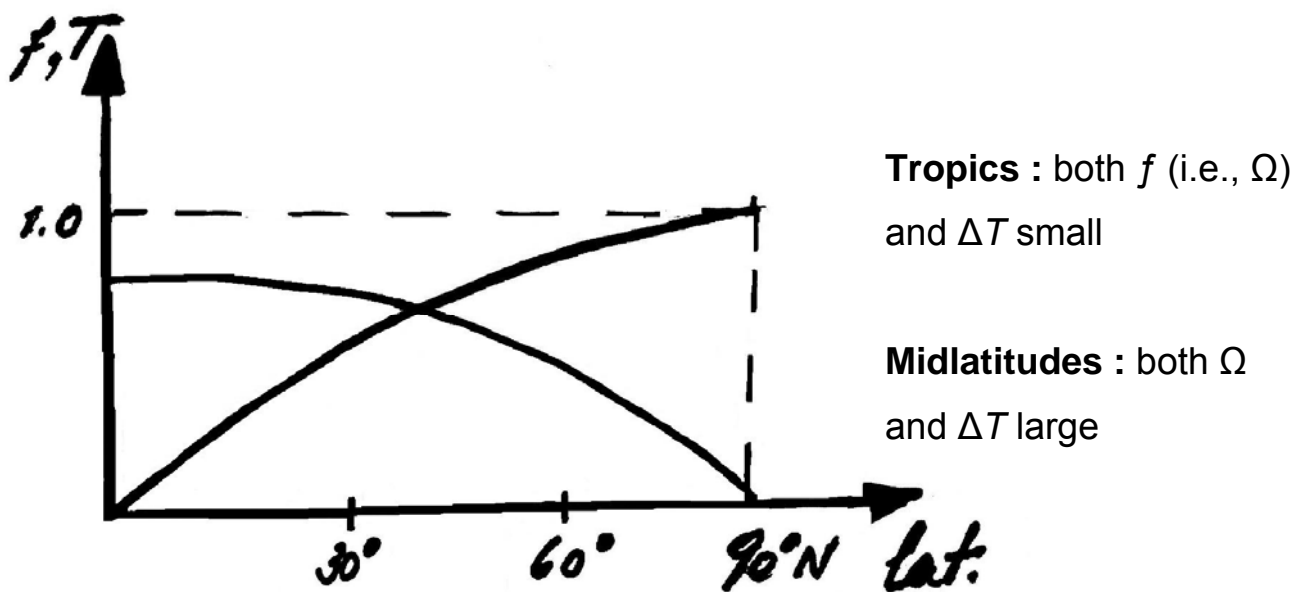
Rotating Convection: An Illustration



Rotating annulus & Earth's atmosphere



Or why doesn't the Hadley cell on Earth extend to the poles, like on Venus ?



Some general references

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